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PHYSIOCHEMICAL PROPERTIES OF AMMONIUM PENTA BORATE HEXAHYDRATESINGLE CRYSTAL: AN EFFICIENT MATERIAL FOR OPTICAL

LIMITINGAPPLICATIONS

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Abstract

AmmoniumPentaBorateHexahydrate(AMPBH),aninorganicnonlinearoptical(NLO)material was

grown by the process of slow evaporation technique (SET) in aqueous solution. Thesynthesized

(AMPBH) was characterized by single crystal X- ray diffraction study to affirm the

structure(monoclinic), lattice parameters and space group P21/n. The functional groups in the material

has been analyzed by FT-IR and FT-RAMAN studies in the range 4000 and 500 cm⁻¹. The optical

absorption

wasdeterminedfromtherecordedUV-

Visiblespectraandenergygapwasevaluatedandthevalueis5.56eV.Thestudyonhardnessinfersthecrystalassoft

materialcategory. The dielectric nature was examined and Jonscher's law confirms the conductivity mechanism

withs=0.984.Laserdamagethreshold study using an Nd;YAG laser (1064nm) having pulse width 10 ns

also carried Theluminescencenatureofcrystalswasinvestigatedintherange300was out.

800nmandconfirmedgreenandvioletfluorescenceemissionspectra. Temperature of decomposition and therma

lstabilitywereinvestigatedusingTGandDTAthermaltechniques.UsingGaussianbeamofHe-

Nelaserofwavelength, the Z-scanmeasurement reveals negative nonlinearity i.e., self-

defocussing, saturation absorption behaviour and thereby the third order nonlinear susceptibility χ(3) were

alsoevaluated.

Keywords: Nonlinearoptical crystal; Borate; Luminescence; Z-Scan;

1. Introduction

The synthesis of nonlinear optical materials to fulfill their contribution in the domain of lasertechnology, optical data storage, optical communication and high speed information technology isgreatlyneedful[1,2].

Frequency conversion materials have the stability to break the restriction in optical spectrum inordertobe utilized innonlinearoptical applications. The desired requirements of the NLO materials rely on large transparency range, high nonlinear susceptibility, fast nonlinear response and higher degree of laser damage threshold (LDT)[3]. Now, as the surge of polyborate salts occupying a vast position in developed applications involving fluorescence, piezoelectric, nonlinear optical (NLO) and porous materials [4-5], the versatility of borates is greatly attractive.

Inorganic borate based crystals are found to possess high transparency in the UV regionand speciality of exhibiting different structures [3] emanating from special bonding properties of boronatom which coordinate with three or four oxygen atoms resulting [BO₃] or [BO₄] groups (AMPBH).Many authors reported borate based materials; Photoinduced NLO effects in doped δ-BiB₃O₆ crystals byMajchrowski et al [6]. Structure-property relationships of inorganic NLO crystal by Chang et al [7] inearlier studies in literature paved way to new series of borate based crystals [8-10]. Also Cook and

JaffeanalysedthepiezoelectricandelectricpropertiesofAmmoniumpentaborate(AB5)[11].C.Chenexplained that the anionic group [B₅O₆ (OH₄)] or [B₅O₈ (OH₂)]³⁻ is responsible for NLO susceptibility ofAB5[12].Alsotheadmiringphysicalpropertiespossessedbyboratesnamelypiezoelectricity,birefringencean despeciallynonlinearopticalnaturearesignificantlyacknowledgedowingtomultifunctionofborateattributedto boro-oxygen(anionicgroups)configurations.Researchersaremeeting a great challenge in design of materials governing the borates in tuning the assembling of theanionic groups. Borates which hold on of both water molecules in it which has poly anions in its crystalstructure.

Borate based NLO materials namely Potassium Borate, Potassium Pent borate (KPB), SodiumTetra borate, Strontium Tetra borate (SrB₄o₇), Lanthanum Calcium borate (La₂CaB₁₀O₁₉), and Bariumstrontiumboratehasbeenextensivelystudied[13]showingstability,higherlaserdamagethreshold

Journal of Vibration Engineering(1004-4523) | Volume 23 Issue 10 2023 | www.jove.science with enhance optical feature and transparent nature in UV region, provides greater difference in theelectro negativity of B and O atoms which provide the basiccharacteristics of borates [14]. A fewboratebasedcandidatewasalsoinvestigatedforfourthandfifthharmonicgenerationusingNeodymium-YAGlaser.

Moreover, Boron and its compounds have a broad range of applications in industry. Especially,boric acid is a well-known medium and its potential rich chemistry has recently been discovered. InBorate materials, the boron atom typically joined with three or four oxygen atoms form [BO₃] or [BO₄]unitsandtheseunitssharecommoncornersprovidingvariousbuildingblockswithdifferentboronatomsfo rmingclusterinasupramolecularmagnesoborate[15].MeantimeB₂O₆unitwithBO₄tetrahedral cluster was synthesized under high pressure with more flexibility was also reported and haveled to the diversity of more borate structures where as far as now more than 2000 borates were reported[15].

Borate crystals of Potassium and Ammonium are excellent non-linear optical materials which issuseful in laser techniques [16] also showing lower absorption to develop Ultraviolet NLO materials. Ammonium pentaborate is a suitable solvent for metal oxides, it is utilized in the process of welding, soldering, and brazing fluxes for stainless steel or nonferrous metals [17]. The effect of shock waves toenhance the optical characteristics of Ammonium Pentaborate Hexahydrate material was studied by Anandarajand Jothi [18].

InthepresentworkweharvestedAmmoniumpentaboratehexahydrate,analkalineboratecrystals from the reaction of Ammonium carbonate, water, and boric acid. Present study includes thesynthesis and growth, single crystal diffraction study, optical characterization, FTIR and FT-Raman,Laser damage threshold study, dielectric, hardness and for the first time Fluorescence andthird ordernonlinearcharacterizationofAMPBHisreported,asuitablecandidatefornonlinearopticalapplications.

2. Experimental

2.1 Materialsynthesisandcrystalgrowth of AMPBH.

Ammonium Penta Borate Hexa Hydrate was synthesized from the starting materials Ammonium carbon at eand boricacido fanalytical grade (AR) procured from Merck Chemicals. The

calculated amount of Ammonium carbonate and boric acid were mixed in stoichiometric molar (1:2)ratio in deionized water at normal temperature. After continuous stirring for 9 hours, the homogenous solution was filtered followed by slow evaporation of the solvent to obtain the material. After 20 daysoptically colorless and good transparent crystals were harvested formed after spontaneous nucleation process. The photograph of as grown crystals of varying sizes are depicted in Fig. 1(a,b).

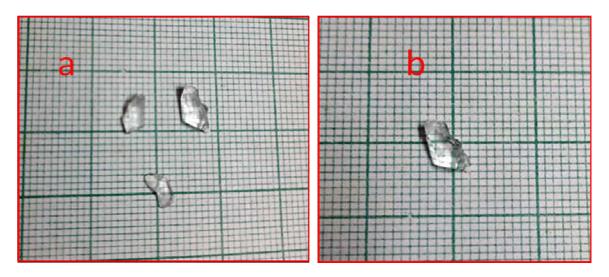


Fig. 1.(a)&(b) As growncrystal of AMPBH

3. Resultsand discussion

3.1 Singlecrystalx-raydiffraction

The grown crystal was given for single crystal X – ray diffraction to analyze the latticeparametersusing BRUKERKAPPAAPEXIICCD single crystal diffraction eteraccompanying MoK α (λ =0.71073Å). X-ray diffraction studies show that the crystal corresponds to the centrosymmetric space group of P21/n holding monoclinic system and the cell parameter values are depicted in Table 1 in accordance with reported work [18] whose structure holds a collection of five Boron with ten Oxyanions and six water molecules. The structure and crystal diffraction pattern of AMPBH is represented in Figure 2.

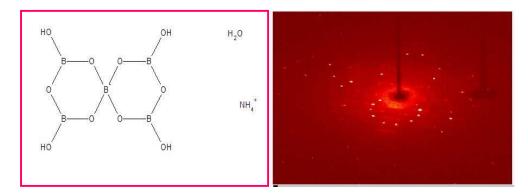


Fig. 2.Structure and Crystal diffractogram of AMPBHTable1.LatticeparametersofAMPBHsinglecrystal

Cell parameters	Presentwork
Spacegroup	P2 ₁ /n
a(Å)	7.18
b (Å)	11.34
c(Å)	7.18
α(°)	90
β(°)	100.16
γ(°)	90
Volume (Å) ³	574
CrystalSystem	MonoclinicP

3.2 FTIRandFT-RamanSpectral Analysis

 $FTIR investigation on Ammonium Penta Borate Hexa Hydrate was determined from Perkin Elmer Spect \\ rophotometer in the range 4000-500 cm^{-1} wavenumber and is depicted in Figure 3 (a) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (b) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (a) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (b) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (b) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in Figure 3 (c) and FT-1000 cm^{-1} wavenumber and is depicted in FT-1000 cm^{-1} wavenumber and i$

RamanspectralanalysiswasperformedusingBrukerRFS27FT-

RamanSpectrometerportrayedinFig.3(b)respectively,wherethebondingofmoleculeinthecrystalelucidatedi tsabsorptionnature. Thepeakoccurredin3125-3378cm⁻¹inIRisdue

toNH₄stretchingandOHasymmetricstretching.TheN-Hstretchingvibrationpeaksobservedat3555cm

¹(IR)andat3551cm⁻¹(FT-

RAMAN) respectively. The O-H (hydroxyl group) stretching of vibrations of water is seen at 3378 cm⁻¹ in IR and the corresponding peak at 3344 cm⁻¹ in Raman spectrum. Also the stretching and bendingmodes of –OH groups are observed at 2459 cm⁻¹ in IR and at 2460 and 2559 cm⁻¹ in FT-Ramanspectrumcounterpart.

Thebandnoticedat1653cm

¹inIRisduetoNH₄asymmetricbendingvibrationsandthecorrespondingpeakobservedat1602cm⁻¹inFT-Raman.Thepeakat1399cm⁻¹inIRcanbeattributedto the asymmetric stretching of the BO₃ groups and the counterpart observed at 1396 cm⁻¹ in FT-Raman.Theborate crystals exhibitmolecular vibrationsin thethree IRspectralregions[19].

- (i) in-between 1500 and 1200 cm⁻¹ pertaining to B-Ostretching of trigonal B-Oin BO₃.
- (ii) The B-O bending vibrations of borate crystal have the absorption bands in the frequency region782-1400cm⁻¹.
- (iii) Bendingvibrations ofborateappearin-between 800and 650cm⁻¹.

The absorption bands observed at 1354, 1248, 1195, 1102, 1025, 694 and 592 cm⁻¹ in IR and peaks observed in FT-Raman spectrum counterpart at 1352, 1240, 1130, 1003, 865, 710 and 616 cm⁻¹

Table2.AssignmentsforFTIRandFT-RAMAN

FT-IR(cm ⁻¹)	FT-Raman(cm ⁻¹)	Assignments
3378	3382	O-HAsymmetricstretching
2459	2460,2559	Stretchingandbending modes of -OH groups
1653	1602	NH ₄ Asymmetricbendingvibrations
1399	1396	AsymmetricstretchingofB-O
1354		B-OAsymmetricstretching
1195,	1205	B-O-Hbendingmode
1248	1240	CH ₂ Torsion
1102	1130	InplanebendingB-O-H
1025	1003	B-OTerminalasymmetricstretching
922	865	B-ORingstretching
		Page No. 6

Page No: 6

¹respectivelycorrespondstoasymmetricandsymmetricstretchingB-Ovibrationsaredetailed inTable 2.

Journal of Vibration Engineering(1004-4523) | Volume 23 Issue 10 2023 | www.jove.science

782		B-ORingstretching
694	710	O-B-ORing symmetricbending
484		O-B-ORing bending
507		O-B-ORing bending
592	562	B-Onetworkin pentaborategroup
592	616	O-B-OTerminal bending

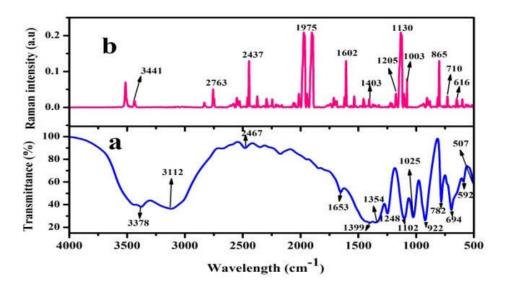


Fig.3.(a) FTIR(b) FT-Ramanprofileof AMPBH

3.3. UV-VisibleStudies

ByemployingPerken-ElmerLambda35UV-Visspectrometeroperatedintherange200-700nm, the transmittance nature of the crystal was analyzed which stands as a key requirement in the field ofphotonics and optoelectronics[19]. The UV-Visible spectra depicted in Fig.4(a) evinces that AMPBHcrystal possess low absorbance and high transmittance in the complete visible area having lower cut-offwavelengthat252 nm is the enticingfactorfor the materialspossessing NLOproperties.

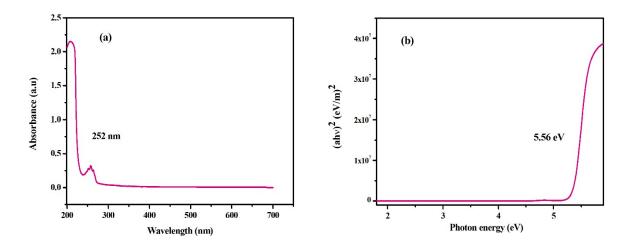


Fig.4.(a,b)UV-visiblespectrum, Tauc'splotforAMPBH

3.3.1 AbsorptionedgeandEnergygapdetermination

 $The calculated absorption values (A) was utilized to evaluate the absorption coefficient (\alpha) utilising the following equation:$

$$\alpha = \frac{2.3026\log(1/T)}{t}$$
 [1]

Where T, timplies the transmittance, tdefines thickness of the crystal. The absorption coefficient (α)was evaluatedby,

$$h\nu\alpha = A(h\nu - E_g)^{1/2}$$
 [2]

Here A is a constant, E_g - energy gap, h infers Plank's constant, υ the frequency of the incident photon. The energy gap of AMPBH crystal was calculated from linear part of the tauc's graph plotted between $(h\upsilon\alpha)^2$ and h υ as depicted in Figure 4(b). The graph exhibits the realization of direct endorsed transition for the titled material in the range 5.0-5.6 eV. From extrapolating the straight line down to $(\alpha h\upsilon)^2 = 0$, direct electronic band gap energy $E_g = 5.56 eV$ was found. Hence, AMPBH has broad optical energy gap can be an approximately a

3.3.2. UrbachEnergy

The absorption coefficient (α) as given by the Urbach equation [20].

$$\alpha(h\nu) = \alpha_0 \exp(h\nu/E_u)$$
 [3]

denote α_o as a constant, E_u pertains urbach energy, h signifies Planck's constant and υ indicates thefrequency of radiation. The plot of photon energy h υ and $ln(\alpha)$ is shownin the profile (Figure(4c)). The crystalline feature of AMPBH is found out from the increase invalue of α with photon energy h υ , depicts high crystallinity in the grown crystal. The slope 3.3316 of the linear part of the graph was calculated from log of the absorption coefficient (α) with photon energy (h υ) that explains that AMPBH has high crystalline nature [19].

The Urbach energy as evaluated from the reciprocal of the slope is 0.3002 eV. Moreover the low(0.3002eV) estimatedUrbach energysignifies minimumdefect in the titleammomiumboratecrystal.

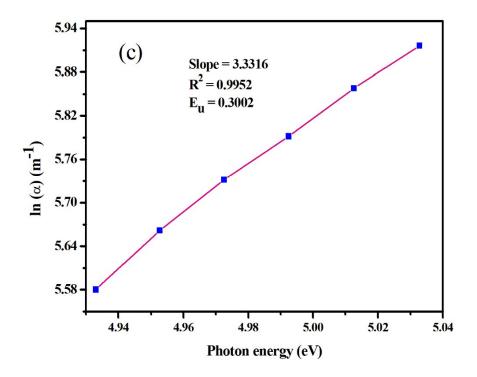


Fig.4c.Plot ofln(α)withphotonenergyforAMPBH

3.3.2 Optical constant determination

Using the spectra of absorption (Fig. 4a), the demanding optical parameters namely transmittance, absorbance, reflectance, absorption coefficient and energy gap were valuated are essentialwhen dealing in the field of processing, tuning, calibrating and design of technology based devices. Inclusive determination of opticalconstants is needed as they promote bringing the conceptregardingtheopticalfeatureofthegrownmaterial. Thus the influence of extinction coefficient (K) and refl ectance(R) of crystal has been studied using the computed transmittance data. The inner efficiency of the appliance too confide on the domains such asabsorption coefficient and extinction coefficient (K)showingpartiallossoflightperunitdistanceinacollectivemedium. The extinction coefficient is evaluatedusingtheformula:

$$K = \frac{\lambda \alpha}{4\pi}$$
 [4]

The graph between extinction coefficient and twain reflectance with photon energy predicts that reduction of extinction coefficient with energy [19] as shown in Fig. 5(a,b). The reflectance (R) in terms of optical absorption coefficient shown below as

$$R=1\pm\frac{\sqrt{1-\exp(-\alpha t)+\exp(\alpha t)}}{1+\exp(-\alpha t)}$$
 [5]

Fig. 6(a,b) depicts the twain reflectance and extinction coefficient confide on absorption coefficient. It is understood that as the values of R and K change with photon energy and also depend on absorption coefficient α , the alteration of optical constants notably the α and E_g contribute significantly towards the fabrication of optical devices [21].

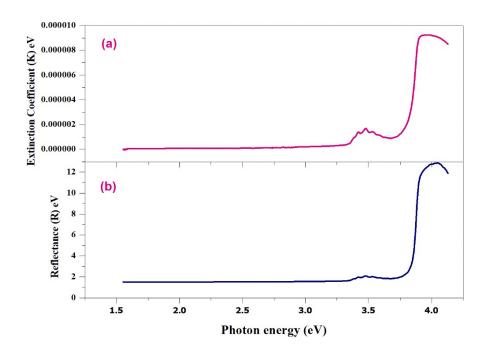
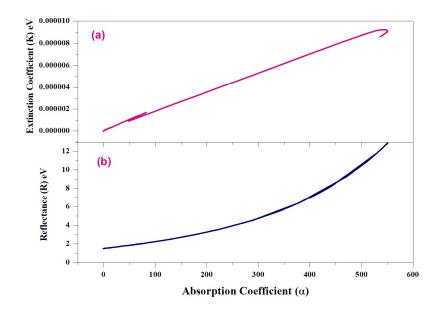


Fig.5(a,,b)Extinctioncoefficientandtwainreflectance(R)dependenceonphotonenergy(hv)



 $Fig.6(a,b)Extinctioncoefficient and twain reflectance(R) dependence on absorption coefficient(\alpha)$

${\bf 3.4.}\ ThermoGravimetric and Differential Thermal Analysis (TG/DTA)$

Thermogravimetricanalysis(TGA)andDifferentialThermalanalysis(DTA)standsasaprincipaltechn iquestocharacterizethethermalstabilityofthegrownmaterial.ForAMBPHthetemperature withstanding point was explored by SII NANO TECHNOLOGY (MODEL TG/DTA 6200)operated between the temperature 20-800 C at a heating rate of 20 C min⁻¹ under nitrogen atmosphere.Initially 2.032 mg was taken for the investigation. The TG/DTA portrayal of AMPBH is shown in Figure 7.

Two main weight losses occur between 125 C-200 C causing liberation of water molecules andNH₃from AMPBH. The endothermic signal in DTA at around 200 C matches with TGA weight losscurve and also at around 25C, the DTA and TGA shows a decline at this temperature (425C) wherethe second weight loss curve ends. Further heating, third weight loss appear at 425 C and continues till 800 C with no further endothermic peaks in DTA and counterpart TGA displaying complete weight loss leading to residue of the material finally at 800°C,

Henceforththe AMPBHcrystal canbesuitable for NLOapplication supto 125C.

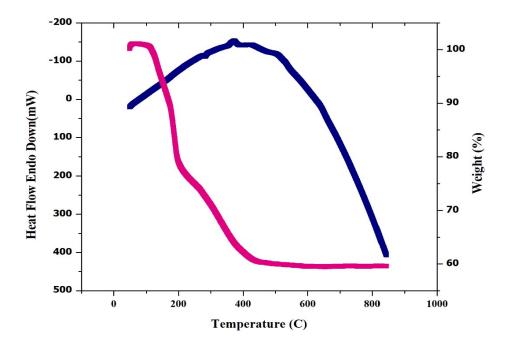


Fig.7.TG/DTAprofileofAMPBH

3.5. Dielectricstudies

The dielectric constant of AMPBH was estimated using the expression,

$$\varepsilon' = \frac{CdA}{\varepsilon_0} \tag{6}$$

Here, d, C represents the thickness and capacitance of the crystal. The change in dielectric constant and dielectric loss with frequency for various temperatures are presented in Fig. 8(a,b) which depicts that dielectric constant and loss decreases increases with frequency. For lesser frequency, dielectric constant gains larger value in view of the presence of all polarization mechanism since low dielectric constant materials has lesser measure of dipoles per unit volume when compared to the crystal possessing high dielectric constant. Hencemeets the criteria in opto-electronic applications.

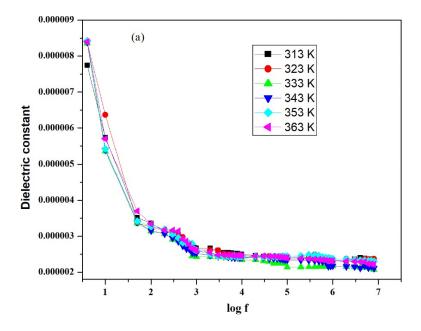


Fig.8.(a) PlotofLogf vsdielectricconstantofAMPBH

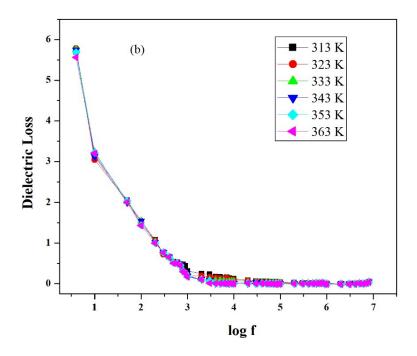


Fig.8.(b) Plotof Logfvsdielectric lossof AMPBH

3.5.1. ActivationEnergy

Activationenergy is obtained from,

$$\sigma = \sigma \quad \exp\left(\frac{E_a}{k}\right) \\
\circ \quad \left(\frac{E_a}{k^B T}\right) |$$
[7]

Here, σ_{ac} —conductivity, T-temperature, E_a - activation energy, K_B - Boltzmann constant (1.38 x 10^{-23} J/K). The graph between log σ_{ac} and 1000/T is showed in Fig.9 from which the activation energy is determined by the expression from the slope in the linear portion,

$$E_a = -Slopex 1000x K_B$$
 [8]

The evaluated values of activation energy was found to be 0.0241, 0.0146, 0.0112 eV at 700 Hz, 3 KHz,20 KHz and at 3 MHz frequencies denotes that inferior values of activation energy, predicts less defects materials pertaining to their usefulness in device fabrications.

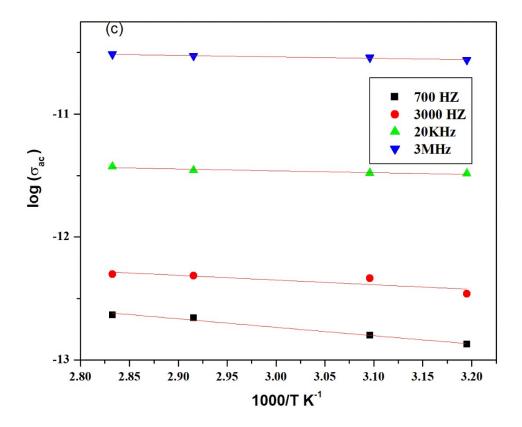


Fig.9.Plotof1000/T versesLog σ_{ac}

3.5.2. MechanismofConduction-Jonscher'spowerlaw

Therelationship between electrical conductivity σ_{ac} and frequency seen invarious solids constituting polymers, glasses and crystals was looked over by Jonscheris well acknowledged by the power law [22],

$$\sigma_{a.c} = \sigma_{d.c} + A\omega^{S}$$
 [9]

, where σ_{dc} infersDC conductivity that concludes the frequency (ω)independent plateau in lowerfrequency, A is a constant that depends on temperature rely on polarisation, and s - power law exponentthat lies in the middle0 < s < 1, that varywith materials. Figure 10 represent the specification of Jonscher's power law explored for the title material whose s = 0.984 established from the slope of the graphleads to the good agreement of the power law. The conduction procedure recognized from the

graphinfersthatACconductivitystep-upalmostathigherfrequencyasaresultofhoppingmechanism [23]observedfromAMPBHcrystal.

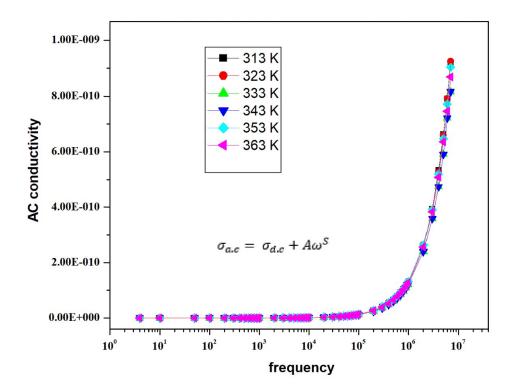


Fig.10. Plotof ACconductivityvs Frequency

3.6. Micro(mechanical) hardnessstudies

The microhardness study confirms the materials capacity in device fabrication by exploring themechanical strength. The strength of the materials depends on various parameters such as lattice energy, Debye temperature, heat of formation and interatomic spacing. The MH- 112 Vicker's microhardnessmeasurement on AMPBH crystal was done at room temperature and load of 25, 50, and 100 g wereapplied.

. The Vicker's hardness which connects the applied load and diagonal of the indentwas evaluated using standard formula

$$H_V = 1.855 (P/d^2)$$
 [10]

Where P implies applied load, dthe mean length of the indenter impression, and 1.8544 is a constant. Figure 11(a) represents the dependence of hardness number (H_v) with applied load extending from 25-

100gforAMPBH,whereH_vincreaseswiththeincreaseintheloadisobservediswellconsideredasthereverseinde ntation sizeeffect[RISE] [24,25].

The Meyer's index coefficient was calculated from Meyer's law [19], which imparts a relation between the load and dindentation diagonal length,

$$P=K^{1}d^{n}$$

$$(or)$$

$$\log P=\log K_{1}+n\log d$$
[12]

Where K₁ signifies constant and 'n' denotes Meyer's index. Work hardening coefficient 'n', ischecked from a plot drawn between log P vs log d as displayed in Figure 11(b). From the slope of thelinearstraightlinethevalueof'n'is3.70.ForthenormalRISEbehavior,theconditionisn<2.Whenn

> 2, it infers reverse RISE behavior. This is good coincidence with the experimental data pointingtowards reverse RISE. As per the statement given by Onitsh [26], work hardening coefficient should bebetween 1 and 1.6 for harder materials and above 1.6 for softer category. Hence, AMPBH can be termedassoftmaterial category.

$$\sigma = \frac{H_{\nu}(0.1)^{n'-2}}{3}$$
 [13]

Where n'=n+2.

Agraphbetweenthevariations of load Pwithyieldstrength σ_y is calculated for AMPBH crystal (Figure 11(c)).

Theelasticstiffnessconstant (C₁₁)isalsofigured out fromWooster'sempirical formula[27],

$$C_{11} = (H_{\nu})^{7/4}$$
 [14]

The elastic stiffness constant recognizes the nature of bonding between neighbouring atoms. A graph isbetween load P vs. stiffness constant C₁₁illustrated in Figure 11 (d) furnish that the stiffness constantincreases with increase of load. The greater the stiffness constant C₁₁lead to the conclusion that the binding forces between theionsarequiteenergetic.

The Knoop's hardness number (H_K) is also determined from,

$$H_K(Kg/mm^2)=14.229\frac{P}{d^2}$$
 [15]

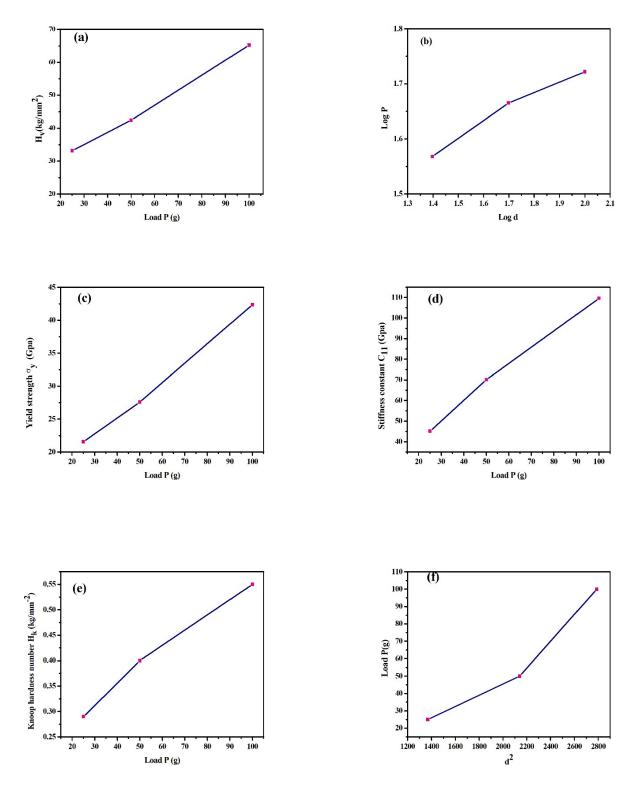
 $From the tabulated values of load Pand H_K (Knoop's hardness number), the graph is plotted (Fig. 11e). \\$ The calculated mechanical parameters are tabulated (Table 3)

3.6.1 Heys-Kendall(HK)Approach

The point of view regarding the study of HK, which initiates the determination of load independent parameter hardness [28], from the relation

$$H_{HK} = 1854.4A_1$$
 [16]

From the hardness analysis AMPBH crystal proves that has sufficient mechanical strength for variousNLOdevicefabrications. The estimatedmechanical parameters are listed in the Table 3.



 $Fig. 11 (a,b,c,d,e) Microhardness \ , \ Meyers \ plot, Yield \ strength, Stiffness \ constant, Knoop's$ $hardness number, load \ Pvs. \ d^2$

Table3:ComputedMechanicalparameters

load P(g)	n	H _v (kgmm ⁻²)	σ _y (G Pa)	H _K (Kg/mm ²)	C ₁₁ (M Pa)
25	3.70	33.18	21.56	25.46	4.495
50	3.70	42.44	27.58	32.56	6.915
100	3.70	65.18	42.35	51.03	14.652

Table4: CalculatedHKconstantW, A₁, andH_{HK}

HKconstant	Results
ResistancePressure(W)	-44.0682(g)
Loadindependent constant(A ₁)	$0.0482(g/\mu m^2)$
Correctedloadindependenthardness(H _{HK})	89.38208(g/μm²)

3.7. PhotoluminescenceStudies(PL).

PLspectroscopyplaysavitalroleinexploringtheelectronicstructureandassessingthetransparencyofth esolidanddiscoveringtheenergystatesexistingbetweenvalencebandandconductionbandresponsibleforradiat iverecombination[28]. The photoluminescences pectrum observed in the range 300-

800nmisdelineatedforAMPBHshowingroomtemperaturephosphorescencewithanexcitationwavelengthof2 53nmusingPERKINELMERLS45Spectrophotometer (Figure 12). The PL spetra portrays one green fluorescence peak at 506 nm (2.45eV), one cyan(blue) emission peak at 463 nm (2.68 eV)and one violet fluorescence peak at 384nm(3.23 eV) which emerges as a result of disappearance of self-trapped excitons linkedwith the moleculartransitions within theBO $^-$ group[19] makes AMPBH suitable in optical data storage devices[29]. Theblue phosphorescence spectra emerges as a result of high energy photons and the generation of greenphosphorescence due to low energy photons [30]. In ammonium boratecrystalsB-O provides the π -bond orbitals and or π^* anti-bond orbitals in phosphorescent excited statesthat paves way to synthesizeadvancedclass ofphosphorescentmaterials applicable in tunablelasers[30].

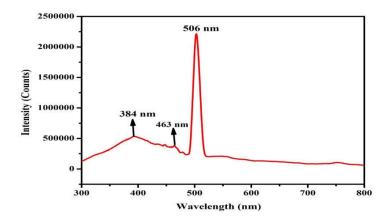


Fig.12.EmissionspectrumofAMPBH

3.8. LaserDamageThreshold(LDT)Study

LDTisusuallyconnected toriseintemperaturein thematerials leadingto strain-inducedfracture evolved by electron avalanche, multi-photon absorption and ionization processes [31]. LDTmeasurement was performed by Q-switched Nd:YAG laser source with primitive wavelength (1064 nm,6ns,10Hz). Multipleshot LDT equivalent of title crystal was evaluated from power density expression:

Powerdensity(
$$P_d$$
)= $E/\tau \pi r^2 (GW/cm^2)$ [17]

[WhereE,\tau,rrepresentsinputenergy(mJ),pulsewidth(ns),radiusofthelaserbeamspot(mm)]is 10.59GW/cm², establishedthematerialssuitability forhigh powerlaserapplications.

3.9. Z-Scanmeasurements

Highly sensitive and a standard method for determining nonlinear optical properties of materials such as solids, liquid solutions and thin films like absolute magnitude and sign of nonlinear refractive index(n₂), nonlinear absorption(NLA) and nonlinear refraction(NLR) was developed forwide applications in the field of optical imaging, multi-photon polymerization including optical switching is the method of Z-scanint roduced by Sheik-

Bahaeetal[32].Heretheclosed(Fig.13a)andopenaperture(Fig.13b)configurations are discussed connected to the emeasurable quantities nonlinear absorption (NLA) and nonlinear refraction (NLR) associated with the imaginary and real part of the third order nonlinear susceptibility $\chi(3)$, provides information about the properties of the material. In the Z-scantechnique used to characterize the AMPBH, He-Nelaserbeam (Gaussian beam) of intensity $5 \, \mathrm{mW}(\lambda=632.8 \, \mathrm{nm})$ was employed as a source, with beam diameter 0.5 mm and a convex lens of focal length

30 mm through which a Gaussian beam was passed. In this procedure, the Z-scan depends in the the the the theoretical transmitted in the spatial and temporal profile of input beam passing through the sample. Total transmitted intensity with respect to the radiations at the sample which is being translated is all due to multi-photon absorption may be limited to two-photon absorption. The variation in far field transmittance beam intensity was measured through the closed aperture using a digital power meter (Field Master GS-coherent). The spatial variation of refractive index caused due to the temperature distribution on surface of the crystal leads to phase distortion in incident beam enables to evaluate the phase-shift ($|\Delta \phi|$) in transmitted beam by the relation,

$$\Delta T_{P-V} = 0.406(1-S)^{0.25} \Delta \Phi \qquad_{0}$$

 $Where (\Delta T_{P-V}) denotes variance between peak and valley transmission in closed aperture assessment and \\ Sisthelinear transmittance aperture and is approximated by,$

$$S=1-\exp\left|\frac{\left(-2r^2\right)}{\left(\omega_a^2\right)}\right|$$
 [19]

 r_a is the radius of the aperture and ω_a is the beam radius and spot diameter at

theaperturesubsequently. The nonlinear refractive index (n₂) can be estimated by using the relation:

$$n_{\overline{2}} = \frac{\Delta \Phi_0}{K I_0 L_{eff}}$$
 [20]

Here $|\phi|$ represents on axis phase shift, where I_0 is the on-axis irradiance at focus (Z=0) (I_0 , the intensity of the laser beam at focus (Z=0) and L_{eff} is the operative thickness

of the sample canbe determinedby

$$L_{eff} = [1 - \exp(-\alpha L)]/\alpha$$
 [21]

Ldenotes the sample length, α represents the linear absorption coefficient, and krepresents the wavenumber ($k = 2\pi/\lambda$).

Inopenaperture(Fig. 13b), the nonlinear transmission of sample without aperture was tallied.

This admits us to resolve the nonlinear absorption coefficient β from open aperture signature,

$$\beta = \frac{2\sqrt[3]{\Delta}T}{I_0 L_{eff}}$$
 [22]

Wherefore $\Delta T=1-T_v$, which that T_v specifies valley transmission value in open aperture determination. This parameter discloses negative value for saturation absorption and positive for two photon absorption and hence an anadvantage to be utilized for optical power limiting process and signal processing devices [28,33].

Thereal andimaginary parts of the thirdorder nonlinear optical susceptibility (χ^3) were established using the formula:

$$\operatorname{Re}\chi^{(3)}(esu) = \frac{10^{-4} (\varepsilon C_0^2 n^2 n)(c m^2)}{\pi} |_{W}$$
 [23]

$$\operatorname{Im}\chi^{(3)}(esu) = \frac{10^{-2}(\varepsilon G^2 n\lambda \beta_0)(cm^2)}{4\pi^2}$$
 [24]

Where, ε_0 is the vacuum permittivity (8.8518x10⁻¹²Fm⁻¹) and λ is the wavelength of laser. The third order nonlinear optical susceptibility (χ^3) of the crystal is figured out from the relation:

$$\left|\chi^{(3)}\right| = \sqrt{(\text{Re}(\chi^{(3)}))^2 + (\text{Im}(\chi^{(3)}))^2}$$
 [25]

The third order nonlinear parameters values checked from data's and formulas were tabulated in Table 5 for AMPBH.

In normalized transmittance curvature collected from the closed aperture (Fig.13a) Z scan data, we can able to look that peak is followed by valley suggesting that the sign of refraction nonlinearity isnegative i.e., self-defocussing. The self-defocussing effect turns out as a result of alteration in refractive index dependence on temperature and hence this property is enforced in the shielding of optical sensornamely night vision devices where IR night-vision couples with IR illumination having spectral range 700-1000nm [29,22].

The nonlinear absorption coefficient (β) value is 0.0209 x 10⁻⁴cm²W⁻¹ and its positive valuestipulates that two photon absorption have occured [28]. The nonlinear optical properties of AMPBHmake it as a capable possible material in the field of optical limiting application. The ratio of closed toopen normaltransmittance function is portrayed in the Fig. 13(c) and Fig. 13(d).

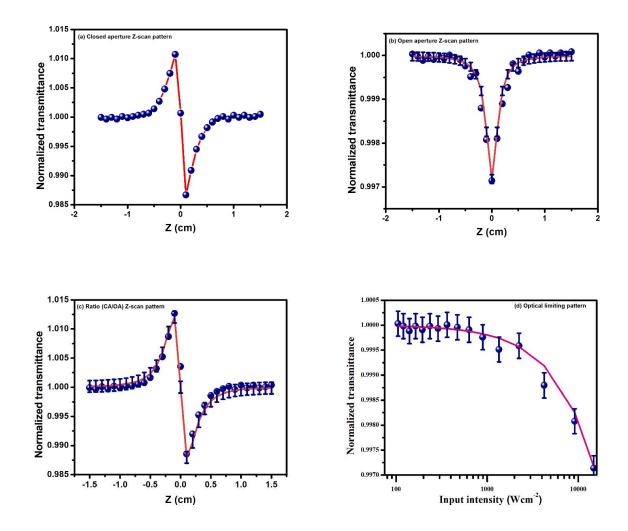


Fig.13.(a,b,c,d) Closed apertureprofile,Openapertureprofile,Ratioofclosedtoopen z-scantraces,
Normalizedtransmittance versusInputIntensity of AMPBH

Table 5: Obtained nonlinear optical parameters from Z-scanme a surements of AMPBH

Journal of Vibration Engineering(1004-4523) | Volume 23 Issue 10 2023 | www.jove.science

Third-ordernonlinearproperties	Measured values	
Nonlinear refractive index	$3.726 \times 10^{-10} \text{cm}^2 \text{W}^{-1}$	
(n ₂)Nonlinear absorption coefficient	$0.0209 \times 10^{-4} \text{cm}^2 \text{W}^{-1}$	
$(β)$ Real susceptibility $(χ^3)$	1.739×10 ⁻⁸ esu	
Imaginary susceptibility	4.129×10 ⁻⁸ esu	
(χ^3) Absolutesusceptibility (χ^3)	4.480×10 ⁻⁸ esu	

4. Conclusion

The alkaline borate symmetric AMPBH crystals were grown and harvested in three weeks byslow evaporation method. The lattice parameters of AMPBH was known from single crystal x-raydiffraction study and found to possess monoclinic system with centrosymmetric space group. FTIR andFT-Raman spectral study established the functional groups comprising the crystal. The evaluated opticalparameters from UV-Vis spectral data clearly suggest the alkaline borate material could be utilized inoptoelectronics and fabrication of devices. AMPBH is thermally stable up to 125 °C which was revealed from TG/DTA analyses. The low dielectric constant and dielectric loss of AMPBH in the high frequency region shows good optical quality. PL spectral profile exhibit sharp green fluorescent peaks at 506 nm,463 nm and at 384 nm corresponding to green and violet fluorescence spectra. The microhardness Vickerstest concludes that the title alkaline borate crystal falls undersoft material category. Theno nlinear refractive index (n_2), absorption coefficient (β) and third-order nonlinear susceptibility χ^3 , (4.480×10⁻⁸ esu) calculated from Z-scan along with optical and spectral analyses concludes that the alkaline borate crystal AMPBH could serve as an admirable candidate for opto-electronic, photonic and nonlinear optical applications.

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